## HOLOGRAPHIC STUDIES OF CELLO VIBRATIONS

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#### 1. INTRODUCTION

Acoustical research on bowed string instruments has traditionally been concentrated on the violin. In spite of being a most important solo and orchestral instrument, the cello has received relatively little attention from musical acoustics researchers. In this paper, we report on studies of the vibrational motion of a cello, using TV holography.

Normal modes are associated with structural resonances. The normal modes of vibration or eigenmodes of a cello are determined by the coupled motion of the top plate, back plate, enclosed air, ribs, neck, fingerboard, etc. Thus one cannot properly label modes as "air modes," "plate modes," or "body modes," although it is tempting to do so when the motion of one part of the structure appears to dominate. We compare the vibrational motion when it is excited by a force applied to the bridge, an oscillating sound pressure applied internally, and the sound field from a loudspeaker, and when it is supported in various ways. From these observations, we attempt to describe the normal modes of vibration of the instrument.

Our designation of the modes follows that used in labeling modes of the violin [1], even though this sometimes results in a different order of mode frequencies in the cello. Perhaps when modal analyses of other cellos are reported, this labeling will change in such a way as to better describe cello modes, as compared with violin modes. The similarity between modes observed in the two instruments is striking, however.

#### 2 SOUND SPECTRA

Room-averaged sound spectra were recorded when the cello was suspended freely by rubber bands, held in normal playing position, and lightly clamped at the center of the neck, the end pin, and at neck/corpus joint (a condition which we will hereafter refer to as "clamped"). The similarity in the peaks in these spectra, shown in Fig. 1, identifies frequencies at which the instrument radiates efficiently.

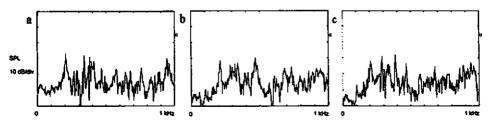


Fig. 1. Room-averaged sound spectra of a cello: a) freely supported on rubber bands; b) hand-held in playing position; c) lightly clamped, as described above.

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## 3. CELLO CAVITY MODES .

To help understand the normal modes of vibration, we studied the air cavity modes of the cello with the top, back, and ribs held nearly stationary by means of sand bags. An oscillating sound pressure was introduced internally by means of a hom compression driver and a length of rubber tubing lagged with cotton to damp acoustic resonances of the tube and driver. Two small microphones (approximately 5 mm in diameter) were used to probe the internal sound field. What we call the (m,n) cavity mode has m transverse nodal planes and n longitudinal nodal planes.

The (0,0) or Helmholtz mode was observed at 104 Hz. The modes characterized by longitudinal motion of the internal air occurred at 226 Hz (0,1), 547 Hz (0,2), 742 Hz (0,3), and 1016 (0,4). The modes characterized by transverse air motion in the lower bout were observed at 496 Hz (1,0) and 893 Hz (2,0). The (1,1) mode, characterized by transverse air motion in both bouts (but stronger in the upper bout), was observed at 609 Hz. Other modes are the (1,2) at 820 Hz the (1,3) at 980 Hz, and the (1,4) at 1129 Hz.

## 4. MODAL ANALYSIS USING HOLOGRAPHIC INTERFEROMETRY

Holographic interferometry offers by far the best spatial resolution of normal modes and operating deflection shapes. Whereas experimental modal testing, and various procedures for mechanical or optical scanning may look at the motion at several hundred (or even thousand) selected points, optical holography looks at almost an unlimited number of points [2]. On the other hand, it is more difficult (but certainly not impossible) to determine the phase of the motion at every point of interest. Furthermore, recording holograms on film tends to be rather time consuming. Electronic TV holography, on the other hand, offers one the opportunity to observe structural motion in real time, and a fast, convenient way to record operational deflection shapes and to determine the normal modes [3].

An optical system for TV holography is shown in Fig. 2. A beam splitter BS divides the laser light to produce a reference and an object beam. The reference beam illuminates the CCD camera via a phase-stepping mirror PS and an optical fiber, while the object beam is reflected by mirror PM so that it illuminates the object to be studied. Reflected light from the object reaches the CCD camera, where it interferes with the reference beam to produce the holographic image. The speckle-averaging mechanism SAM in the object beam alters the illumination angle in small steps in order to reduce laser speckle noise in the interferograms.

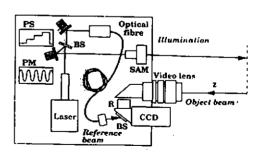


Fig. 2. Optical system for TV holography

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## 5. NORMAL MODES OF THE CELLO

The  $C_1$  mode at 57 Hz is a one-dimensional bending or beam mode. This mode radiates almost no sound, but it contributes to the "feel" of the instrument. The first radiating mode is the  $A_0$  or f-hole mode at 102 Hz, which is shown in Fig. 3. This mode can be excited by a horizontal force on the bridge, an oscillating internal pressure, or by the sound field of a loudspeaker. In other cellos, the  $A_0$  mode has been observed at frequencies as low as 82 Hz [4]. The reason for this rather large difference is not clear, since cello bodies are all rather similar in size and shape.

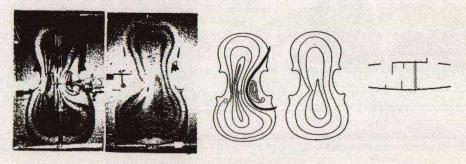


Fig. 3. A<sub>0</sub> mode in the top and back plates. The diagram at the right is a slice view in the bridge plane showing the amplitude and direction of motion at various points.

The  $T_1$  mode, shown in Fig. 4, is observed at 144 Hz. The top plate motion is quite similar to that observed in the  $T_1$  violin mode, but the motion of the back is somewhat different, and it appears to radiate less efficiently than in the violin.

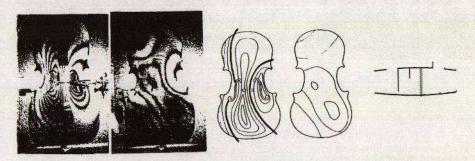


Fig. 4. T, mode at 144 Hz excited by a transverse bridge force. The diagram at the right is a slice view in the bridge plane showing the amplitude and direction of motion at various points.

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The  $C_2$  mode, characterized by a torsional or twisting motion, as shown in Fig. 5, is observed at 170 Hz. Unlike the violin, this mode occurs at a higher frequency than the  $T_1$  mode.

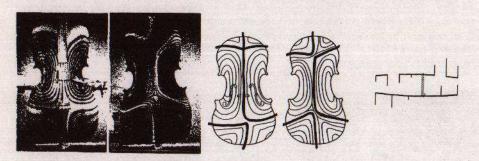


Fig. 5. C<sub>2</sub> mode at 170 Hz. The diagram at the right is a slice view in the bridge plane showing the amplitude and direction of motion at various points.

A mode with ring-shaped nodes in the top and back plates, similar to the C<sub>4</sub> mode in a violin, is found at 195 Hz. This mode is shown in Fig. 6.

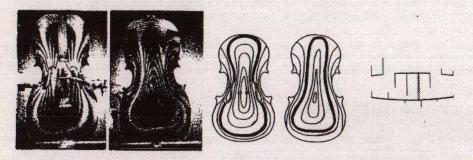


Fig. 6. C<sub>4</sub> mode with ring-shaped nodes in the top and back plates. The diagram at the right is a slice view in the bridge plane showing the amplitude and direction of motion at various points.

In the  $A_1$  mode at 203 Hz, the internal air moves longitudinally between the upper and lower bouts, as in the (0,1) cavity mode. The plates more or less pivot on a nodal line near the bridge as shown in Fig. 7.

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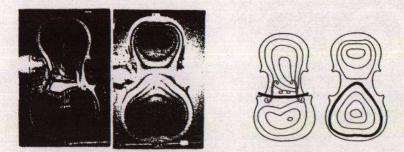


Fig. 7. A<sub>1</sub> mode excited with oscillating internal pressure. The internal air moves longitudinally between the upper and lower bouts, as in the (0,1) cavity mode.

The strongest radiating mode in this range, as in the violin, is the  $C_3$  mode at 211 Hz, shown in Fig. 8, at a frequency above that of the  $C_4$  mode. Note that the back vibrates more strongly than the top plate. In an early 19th century cello, Askenfelt [5] observed the  $T_1$ ,  $C_3$ , and  $C_4$  modes clustered closely together just below 200 Hz.

With the A-string damped, the C<sub>3</sub> mode occurs at 211 Hz. When the A-string is not damped, however, a pair of strong modes with identical plate motion are observed at 207 and 220 Hz. In between, at 219 Hz, we observe a sharp minimum in response, which indicates that the vibrating string opposes the horizontal force applied to the bridge.

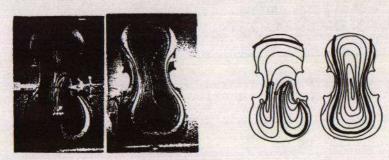


Fig. 8. C<sub>3</sub> mode at 211 Hz (A-string damped).

The  $A_2$  mode at 302 Hz is characterized by longitudinal (0,1) motion of the internal air, similar to the  $A_1$  mode but with the opposite phase. The frequencies of the  $A_2$  and  $A_1$  modes in the cello have the ratio 1.5, compared to a ratio of 1.75 (or 7:4) in the violin. The plate motion, shown in Fig. 9, is quite similar to that of the  $A_1$  mode.

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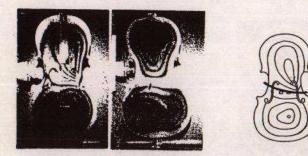
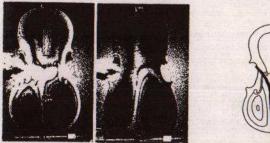


Fig. 9. A<sub>2</sub> mode excited with oscillating internal pressure. The internal air moves longitudinally between the upper and lower bouts, as in the A<sub>1</sub> mode but with the opposite phase.

In the  $A_3$  mode at 277 Hz, the internal air appears to move transversely in the lower bout, as in the (1,0) cavity mode. The plates vibrate in accord with the changing internal pressure.



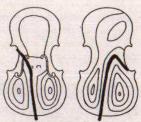


Fig. 10. A<sub>3</sub> mode at 277 Hz excited with oscillating internal pressure. The internal air moves transversely in the lower bout, and the plates vibrate accordingly.

## 6. DISCUSSION

In the table below, we have tabulated the modal frequencies observed in this cello, along with those we were able to find in the published literature. A comparison is difficult, since modal shapes are not given in most of the publications. Nevertheless, the agreement appears to be satisfactory in most cases. If one interchanges Askenfelt's  $C_4$  and  $C_3$  mode frequencies, the agreement is even better. In at least three of the cellos, there is an octave spacing between the  $A_1$  and  $A_0$  modes, which appears to be the case in most cellos

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[6]. Firth [7] describes a number of modes as top-plate modes (T) or back plate modes (B). Langhoff [8] lists several modes of the neck and tailpiece but only three modes of the body.

Mode	Bynum 57 Hz	Ratio to	violin	Comments	Askenfe	elt.	Langho	ff	Firth 1	Firth 2E	ggers
C₁ A₀	102	0.36				90 Hz	85 Hz		90 Hz	104 Hz	82Hz
T <sub>1</sub>	144	0.32	Stronge	er in top than ba	ck	168	146		165	185	132
C <sub>2</sub>	170	0.42	Torsion	nal motion			163				186?
C.	195	0.29	Ring-sl	naped node		202					
$\mathbf{A}_1$	203	0.43	Longitu	ıdinal air motio	n				188	203	
C <sub>3</sub>	219	0.39				185					218?
$A_3$	277	0.25	Transv	erse air motion							260
A <sub>2</sub>	302	0.37	Longitu	idinal air motio	n						312
								T <sub>2</sub>	290	292	
								T,	380	330	
								$T_4$		440	
								$\mathbf{B}_{1}$	178	205	
								B <sub>2</sub>	295	283	
								В,	320	308	
								B <sub>4</sub>	390	390	
								$\mathbf{B}_{s}$		460	

In the table above, we give the ratios of the cello mode frequencies compared to those observed in the Hutchins violin SUS 295 [1]. The ratio of cello string frequencies to violin string frequencies is 0.33, but their body lengths and widths are more nearly in the ratio of 0.5. The cello is not meant to be a scaled up violin, of course, but the comparisons are interesting. A good discussion of simularity laws is found in Cremer [9].

#### 7. CONCLUSIONS

Holographic interferometry is an accurate and convenient way to do modal analysis of the cello and a valuable aid in determining the normal modes of vibration. Electronic TV holography is especially convenient, since the mode shapes can be displayed essentially in real time.

The modes of a cello are found to be quite similar to the corresponding modes of a violin, although they may appear in a different order in frequency. Modal frequencies in this cello occur at about 0.29 to 0.43 times the corresponding mode frequencies in a violin. Although the cello is certainly not a scaled-up violin, it does belong to the "violin family" of instruments.

## 7. REFERENCES

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