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### POSITION FIXING OF TRANSPONDERS FROM SURFACE MEASUREMENTS

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#### 1. INTRODUCTION

There are many underwater engineering applications in which an array of acoustic transponders is needed on the sea bed [1-5]. If the applications are for accurate position fixing it is necessary to know their exact coordinates. In practice, the transponders may be dropped onto the sea bed from a ship and therefore these coordinates may not be known accurately. With *intelligent* devices it is possible to determine their separations by programming them to interrogate each other [4]. The problem posed here is how to determine the geometry and orientation of an array from surface measurements only, a practical requirement suggested by an offshore company.

The problem is approached by considering an array of three transponders A, B and C, as shown in Fig.1. The requirement is to determine the distances  $d_1$ ,  $d_2$  and  $d_3$  between each pair of transponders and the angle  $\theta$  between the line AB and the direction AN of true or magnetic north.

Attention is limited to the case where the sea bed is a plane, which can be horizontal or slanting, and the sea water is a stratified medium with a known sound velocity profile (SVP). Four situations are considered: (i) flat sea bed, constant SVP, i.e. c(z) = c; (ii) flat sea bed, any SVP; (iii) slanting sea bed, constant SVP; (iv) slanting sea bed, any SVP.

The practical approach is to make measurements at two points on the surface. The measurements needed in all four situations are illustrated in Fig.2. At a surface observation point  $O_1$ , the depth  $h_1$  is measured with a depth sounder and the acoustic propagation times  $t_{A1}$ ,  $t_{B1}$ ,  $t_{C1}$  along rays  $O_1A$ ,  $O_1B$ ,  $O_1C$  are measured by interrogating each transponder in turn. At a second surface observation point  $O_2$ , the depth  $h_2$  and times  $t_{A2}$ ,  $t_{B2}$ ,  $t_{C2}$  are measured. The length  $O_1O_2$  and the angle  $\phi$  between this line and true or magnetic north N are also measured. For situations (3) and (4), additional measurements are required to give the directions of each of the transponders as seen from  $O_1$  and  $O_2$ , that is  $\gamma_{A1}$ ,  $\gamma_{B1}$ ,  $\gamma_{C1}$ ,  $\gamma_{A2}$ ,  $\gamma_{B2}$ ,  $\gamma_{C2}$ .

The key feature of the approach is to project all the construction lines  $O_1A,.....,O_2C,O_1O_2$ , and the angle  $\phi$  onto the plane of the sea bed. The problem is thus reduced to two-dimensional geometry [2], as shown in Fig.3(a). As long as all the projected lines  $P_1A, P_1B, P_1C, P_2A, P_2B, P_2C, P_1P_2$ , and the angle  $\psi$  are determined, calculation of  $d_1$ ,  $d_2$ ,  $d_3$  and  $\theta$  is straightforward. To illustrate the procedure,  $d_1$  and  $\theta$  are calculated. The relevant geometry is shown in Fig.3(b). Applying the cosine theorem to triangles  $AP_1P_2$  and  $BP_1P_2$ ,

$$\cos \angle AP_1P_2 = [P_1A^2 + P_1P_2^2 - P_2A^2] / 2 P_1A.P_1P_2$$
(1)

$$\sin[AP_1P_2] = \sqrt{[1 - \cos^2[AP_1P_2]]}$$
 (2)

$$\cos \angle BP_1P_2 = [P_1B^2 + P_1P_2^2 - P_2B^2] / 2 P_1B.P_1P_2$$
(3)

$$\sin \left( BP_1P_2 = \sqrt{1 - \cos^2 \left( BP_1P_2 \right)} \right) \tag{4}$$

Since  $\angle AP_1B = \angle AP_1P_2 - \angle BP_1P_2$ , then

$$\cos(AP_1B) = \cos(AP_1P_2)\cos(BP_1P_2) + \sin(AP_1P_2)\sin(BP_1P_2)$$
 (5)

Applying the cosine theorem again to AP1B yields

$$d_1 = \sqrt{[P_1A^2 + P_1B^2 - 2 P_1A.P_1B.\cos(AP_1B)]}$$
 (6)

Similarly,

$$d_2 = \sqrt{[P_1B^2 + P_1C^2 - 2 P_1B.P_1C.cos/BP_1C]}$$
 (7)

$$d_3 = \sqrt{(P_1C^2 + P_1A^2 - 2 P_1C.P_1A.\cos/CP_1A)}$$
(8)

where  $\cos \angle BP_1C$  and  $\cos \angle CP_1A$  can be calculated by variable replacement  $(A \rightarrow B \rightarrow C \rightarrow A)$  in equations (1) to (5).

To calculate  $\theta$ , consider  $\triangle AP_1M$  in Fig.3(b), where M is the cross-over point of  $P_1P_2$  with the line drawn from A in the direction of true or magnetic north N. Simple geometric relations give

$$\theta = \angle P_1 AB - (180^\circ - \angle AP_1 P_2 - \psi) \tag{9}$$

where

$$\angle AP_1P_2 = \cos^{-1}[\cos\angle AP_1P_2] \tag{10}$$

$$\angle P_1AB = \cos^{-1}\{[P_1A^2 + AB^2 - P_1B^2] / 2 P_1A.AB\}$$
 (11)

Thus, to locate the transponder positions it is necessary to calculate all the projection lines and angles. In the following sections these are calculated for the situations listed earlier.

# 2. FLAT SEA BED, CONSTANT SOUND VELOCITY PROFILE

If the sound velocity does not vary with depth, i.e. c(z) = constant, each acoustic ray is a straight line with a length equal to the propagation time multiplied by the sound velocity c. For example,  $O_1A = c.t_{A1}$ . For a flat bottom,  $h_1 = h_2 = h$  and all the projected lines  $P_1A$ ,  $P_1B$  and  $P_1C$  are perpendicular to  $O_1P_1$  and  $O_2P_2$ , as shown in Fig.2, i.e.

$$P_1A^2 = (c.t_{A1})^2 - h^2 (12)$$

$$P_1B^2 = (c.t_{B1})^2 - h^2$$
 (13)

$$P_1C^2 = (c.t_{c1})^2 - h^2$$
 (14)

$$P_2A^2 = (c.t_{A2})^2 - h^2 (15)$$

$$P_2B^2 = (c.t_{B2})^2 - h^2 ag{16}$$

$$P_1C^2 = (c.t_{c2})^2 - h^2$$
 (17)

It follows that  $P_1P_2 = O_1O_2$  and  $\psi = \phi$ .

### 3. FLAT SEA BED, ANY SOUND VELOCITY PROFILE

In this situation the acoustic rays are curved lines in general and the first consideration is how to calculate  $P_1A$  from the geometry shown in Fig.4. According to two widely known formulae in the theory of sound transmission [5],  $P_1A$  may be given in terms of a ray parameter 'a', i.e.

$$a = \sin \eta / c(z) \tag{18}$$

$$t_{A1} = \int_{0}^{\infty} dz / c(z) \sqrt{[1 - a_{A1}^{2} c^{2}(z)]}$$
 (19)

$$P_1A = \int_0^1 a_{A1} c(z) dz / \sqrt{[1 - a_{A1}^2 c^2(z)]}$$
 (20)

In (19) and (20),  $a_{A1}$  is the ray parameter for a ray  $O_1A$  having an initial angle  $\eta_0$ . Equation (19) has a general form given by  $f_1(a_{A1}) = 0$ . This is a known non-linear function that can be calculated by numerical methods. Applying a conventional iterative method, e.g. the Newton-Raphson method, a solution can be obtained for  $a_{A1}$ . The approach is to try many values of  $\eta_0$ , then choose the one that leads to a vertical depth  $h_1$  during a specified time  $t_{A1}$ . After  $a_{A1}$  has been solved,  $P_1A$  can be calculated numerically with (20). This method leads to a general procedure for calculating all the projected lines of  $P_1A, \ldots, P_2C$  ( $P_1P_2 = O_1O_2$  and  $\psi = \phi$  as before):

Solve a non-linear equation

$$f_{Ti}(a_{Ti}) \equiv t_{Ti} - \int_{0}^{\infty} dz / c(z) \sqrt{[1 - a_{Ti}^{2} c^{2}(z)]} = 0$$
 (21)

for  $a_{Ti}$ , where T = A, B, C; i = 1, 2.

Calculate P<sub>i</sub>T

$$P_{i}T = \int_{0}^{\infty} a_{Ti} c(z) dz / \sqrt{1 - a_{Ti}^{2} c^{2}(z)}$$
 (22)

### 4. SLANTING SEA BED, CONSTANT SOUND VELOCITY PROFILE

The geometry related to the calculation of P<sub>1</sub>A is shown in Fig. 5, which has the following features:

- The ray  $O_1A$  is a straight line whose length is  $O_1A = c t_{A1}$ ;
- Pour lines on a horizontal plane are drawn from  $P_1$ : the true or magnetic north direction  $P_1N$ , the direction of A as seen from  $O_1$ , the steepest direction (at angle β) and the zero gradient direction of the sea bed;
- The direction of transponder A and the steepest direction have angles  $\gamma_{A1}$  and  $\sigma$  respectively with respect to P<sub>1</sub>N;
- A<sub>h</sub> and A<sub>v</sub> are the projection of A onto the horizontal plane and the vertical line O<sub>1</sub>P<sub>1</sub>.

The following theorem is now used: if a slanting plane has an oblique angle  $\beta$  along its steepest direction, then the oblique angle  $\delta$  in a direction at an angle  $\chi$  to the steepest direction will satisfy

$$tan \delta = tan \beta.cos \chi \tag{23}$$

Applying (23) to the direction of A from P, gives

$$tan(\angle AP_1A_h) = tan \beta.cos(\gamma_{A1} - \alpha)$$
 (24)

$$z_{A1} = h_1 - AA_h = h_1 - x_{A1} \tan \beta .\cos(\gamma_{A1} - \alpha)$$
 (25)

From  $\Delta O_1 AA_v$ ,

$$x_{A1}^2 + z_{A1}^2 = O_1 A^2 = c^2 t_{A1}^2$$
 (26)

Combining (25) and (26) gives the solution of  $x_{A1}$  and  $z_{A1}$ :

$$[1 + \tan^2\beta \cdot \cos^2(\gamma_{A1} - \sigma)] x_{A1}^2 - [2h_1 \tan \beta \cdot \cos(\gamma_{A1} - \sigma)] x_{A1} + [h_1^2 - c^2 t_{A1}^2] = 0$$
 (27)

This equation leads to a solution for  $x_{A1}$ . In  $\Delta P_1AA_h$ ,

$$P_1A = x_{A1} \sec((AP_1A_1) = x_{A1}\sqrt{1 + \tan^2\beta \cdot \cos^2(\gamma_{A1} - \alpha)})$$
 (28)

After algebraic manipulation, this becomes

$$h_{1} \tan \beta .\cos(\gamma_{A1} - \alpha) + \sqrt{\{c^{2} t_{A1}^{2} [1 + \tan^{2}\beta .\cos(\gamma_{A1} - \alpha)] - h_{1}^{2}\}}$$

$$P_{1}A = \sqrt{\{1 + \tan^{2}\beta .\cos^{2}(\gamma_{A1} - \alpha)\}}$$
(29)

 $P_1B$ ,  $P_1C$ ,  $P_2A$ ,  $P_2B$ ,  $P_2C$  can each be calculated with a formula similar to (29) by replacing  $t_{A1}$ ,  $\gamma_{A1}$  by  $t_{B1}$ ,  $\gamma_{B1}$ ....., and  $h_1$  by  $h_2$  where necessary. It is then possible to write a general formula for calculating all six projected lines as

$$P_{i}T = \{h_{i} Q_{Ti} + \sqrt{[c^{2}t_{Ti}^{2}(1 + Q_{Ti}^{2}) - h_{i}^{2}]} / \sqrt{[1 + Q_{Ti}^{2}]}$$
(30)

$$Q_{Ti} = \tan \beta . \cos(\gamma_{Ti} - a) \tag{31}$$

where T = A, B, C and i = 1, 2.

To obtain  $P_1P_2$  and  $\psi$ , Fig. 6 is used, where  $O_1$  and  $O_2$  are the two surface observation points;  $O_1O_3 = O_1O_2$ ;  $P_1$ ,  $P_2$ ,  $P_3$  are the projected points of  $O_1$ ,  $O_2$ ,  $O_3$  respectively onto the sea bed; and  $P_{2h}$  and  $P_{3h}$  are the projected points of  $O_2$  and  $O_3$  onto the horizontal plane which passes through  $P_1$ . It is evident that for the trapezium  $O_1P_1P_2O_2$  and  $\Delta P_2P_1P_{2h}$  the following two equations hold:

$$P_1 P_2^2 = O_1 O_2^2 + (h_1 - h_2)^2$$
(32)

$$tan(\angle P_2 P_1 P_{2h}) = (h_1 - h_2) / O_1 O_2$$
(33)

and when applying (23),

$$tan(\angle P_2 P_1 P_{2h}) = tan \beta.cos(\phi - \sigma)$$
(34)

Combining (33) with (34) gives  $\alpha$  as

$$\alpha = \phi - \cos^{-1} \{ [h_1 - h_2] / [O_1 O_2 \tan \beta] \}$$
 (35)

This calculation shows that it is not necessary to know a beforehand. From  ${}_{\Delta}P_{3}P_{1}P_{3h}$  and  ${}_{\Delta}P_{2}P_{1}P_{2h}$ 

$$P_{3}P_{3h} = P_{1}P_{3h}\tan(\angle P_{3}P_{1}P_{3h}) = O_{1}O_{2}\tan\beta.\cos\alpha$$
 (36)

$$P_1P_2 = P_1P_{2b} \sec(\angle P_2P_1P_{2b}) = O_1O_2\sqrt{1 + \tan^2\beta \cdot \cos^2\alpha}$$
 (37)

$$P_{2}P_{2h} = P_{1}P_{2h} \tan(\angle P_{2}P_{1}P_{2h}) = O_{1}O_{2} \tan \beta \cdot \cos(\phi - \alpha) = h_{1} - h_{2}$$
(38)

$$P_1P_2 = P_1P_{2h} \sec(\angle P_2P_1P_{2h}) = O_1O_2\sqrt{1 + \tan^2\beta \cdot \cos^2(\phi - \sigma)}$$

$$= O_1O_2^2 + [h_1 - h_2]^2$$
 (39)

From trapezium P<sub>3</sub>P<sub>3h</sub>P<sub>2h</sub>P<sub>2</sub>

$$P_2P_3^2 = P_{2h}P_{3h}^2 + (P_2P_{2h} - P_3P_{3h})^2 (40)$$

where  $P_{2h}P_{3h}$  can be calculated using cosine theorem in  ${}_{\Delta}P_{1}P_{2h}P_{3h}$  as

$$P_{2h}P_{3h}^2 = P_1P_{2h}^2 + P_1P_{3h}^2 - 2P_1P_{2h}P_1P_{3h}\cos\phi$$

$$= 2O_1O_2^2 (1 - \cos\phi)$$
(41)

All three sides of  $\Delta P_1 P_2 P_3$  are now known. The cosine theorem again results in

$$\cos \psi = [P_1 P_2^2 + P_2 P_3^2 - P_2 P_3^2] / 2 P_1 P_2 P_2 P_3$$
(42)

After manipulation, the result is

$$\psi = \cos^{-1} \frac{\cos \phi + (h_1 - h_2) \tan \beta .\cos \alpha / O_1 O_2}{\sqrt{[1 + \tan^2 \beta .\cos^2 \alpha]} \sqrt{[1 + (h_1 - h_2)^2 / O_1 O_2]}}$$
(43)

It should be noted that when  $\beta = 0$ ,  $h_1 = h_2$ , the formulae above reduce to (12) through (19), the situation for a flat sea bed.

## 5. SLANTING SEA BED, ANY SOUND VELOCITY PROFILE

Fig. 5 is also applicable to this situation. The only difference is that the ray O<sub>1</sub>A is now a curve; so (25) is still correct but (26) is not. Instead of (25) two acoustic transmission formulae similar

to (19) and (20) must be used. Therefore the three equations for solving  $a_{A1}$ ,  $x_{A1}$  and  $z_{A1}$  are:

$$t_{A1} = \int_{0}^{\infty} dz / c(z) \sqrt{1 - a_{A1}^{2} c^{2}(z)}$$
 (44)

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$$x_{A1} = \int_{0}^{\infty} a_{A1} c(z) dz / \sqrt{[1 - a_{A1}^{2} c^{2}(z)]}$$
 (45)

$$z_{A1} = h_1 - x_{A1} \tan \beta .\cos(\gamma_{A1} - \alpha) \tag{46}$$

Inserting (44) into (45) yields

$$h_1 - z_{A1} = \tan \beta .\cos(\gamma_{A1} - \alpha) \int_0^\infty a_{A1} c(z) dz / \sqrt{[1 - a_{A1}^2 c^2(z)]}$$
 (47)

Equations (44) and (47) have the general form

$$f_1(z_{A1}, a_{A1}) = 0 (48)$$

$$g_1(z_{A1}, a_{A1}) = 0$$
 (49)

where functions  $f_1(.)$  and  $g_1(.)$  are non-linear (but known) and can be calculated numerically. Equations (48) and (49) can be solved by conventional iteration. As long as  $z_{A1}$  and  $a_{A1}$  are available,  $x_{A1}$  can be easily calculated by using (46). The desired projected line  $P_1A$  is therefore obtained as

$$P_1A^2 = x_{A1}^2 + (h_1 - z_{A1})^2 (50)$$

The whole procedure for calculating P<sub>1</sub>A,....,P<sub>2</sub>C in a general form can be expressed as follows:

Solve a set of non-linear equations as

$$f_{Ti}(z_{Ti}, a_{Ti}) \equiv t_{Ti} - \int_{0}^{\infty} dz / c(z) \sqrt{[1 - a_{Ti}^{2} c^{2}(z)]} = 0$$
 (51)

$$g_{Ti}(z_{Ti}, a_{Ti}) = h_i - z_{Ti} - \tan \beta .\cos(\gamma_{Ti} - \alpha) \times \int_0^{\infty} a_{A1} c(z) dz / \sqrt{[1 - a_{Ti}^2 c^2(z)]} = 0$$
 (52)

for  $z_{Ti}$  and  $a_{Ti}$ , where T = A, B, C; i = 1, 2.

Calculate x<sub>τi</sub>

$$\mathbf{x}_{\mathsf{T}i} = [\mathbf{h}_i - \mathbf{z}_{\mathsf{T}i}] / \tan \beta. \cos(\gamma_{\mathsf{T}i} - a) \tag{53}$$

Calculate P<sub>i</sub>T

$$P_{i}T = x_{Ti}^{2} + (h_{i} - z_{Ti})^{2}$$
 (54)

 $P_1P_2$ ,  $\sigma$  and  $\psi$  are still given by (32), (35) and (43).

### 6. SUMMARY

A projection approach has been presented for calculating the geometry of an array of underwater acoustic transponders on the sea bed from two surface measurements only. The algorithms can give the array's configuration, size and orientation with respect to true or magnetic north. Four situations have been considered, those of a flat or slanting sea bed each for either a constant or variable sound velocity profile. The procedure can be divided into three steps:

#### Measurements:

- the acoustic propagation times from two observation points  $O_1$  and  $O_2$  to each transponder A, B and C:  $t_{A1}$ ,  $t_{B1}$ ,  $t_{C1}$ ,  $t_{A2}$ ,  $t_{B2}$ ,  $t_{C2}$  (Fig. 2)
- the sea depth at O<sub>1</sub> and O<sub>2</sub>: h<sub>1</sub>, h<sub>2</sub>
- the length and direction of the line connecting O<sub>1</sub> and O<sub>2</sub>: O<sub>1</sub>O<sub>2</sub>, φ
- (slanting sea bed only) the directions of each transponder seen from  $O_1$  and  $O_2$ :  $\gamma_{A1}$ ,  $\gamma_{B1}$ ,  $\gamma_{C1}$ ,  $\gamma_{A2}$ ,  $\gamma_{B2}$ ,  $\gamma_{C2}$  (Fig. 2); the oblique angle  $\beta$  along the steepest direction must be known, but the steepest direction  $\sigma$  itself does not need to be known (Fig. 5).

### Projection:

The aim is to calculate the projected lines of all the acoustic rays  $O_1A$ ,  $O_1B$ ,  $O_1C$ ,  $O_2A$ ,  $O_2B$ ,  $O_2C$ , as well as  $O_1O_2$ , onto the sea bed (Fig. 2). Also needed is the projected angle of  $\phi$ . The equations for all four situations are summarised as follows:

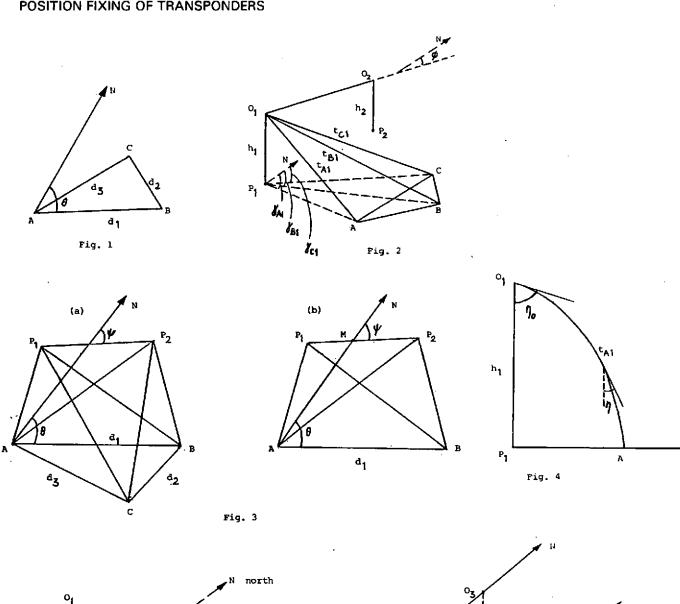
0	Flat sea bed, constant SVP	12, 13, 14, 15, 16, 17
0	Flat sea bed, any SVP	21, 22
0	Slanting sea bed, constant SVP	30, 31, 32, 35, 43
٥	Slanting sea bed, any SVP	<b>51, 52,</b> 53, 54, 32, 35, 43

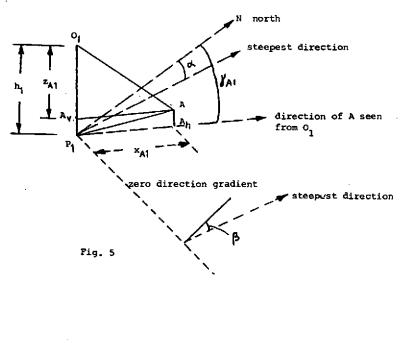
### Location:

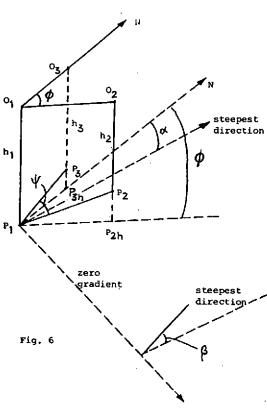
When all the projected lines and angles have been calculated, the whole problem reduces to a plane geometrical one and the algorithms are given by equations (1) through (11), which are common to all four situations. In this paper the calculations have been done for a three-transponder array for the sake of illustration. Because an array of any number of transponders can be considered to be a combination of triangles, the method can easily be extended.

## 7. REFERENCES

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